

## EVALUATION OF THE GAS PRODUCTION POTENTIAL OF CHALLENGING HYDRATE DEPOSITS

George J. Moridis, Matthew T. Reagan, Katie L. Boyle, and Keni Zhang

Lawrence Berkeley National Laboratory  
1 Cyclotron Rd.  
Berkeley, CA 94720, USA  
e-mail: GJMoridis@lbl.gov

### ABSTRACT

We use the TOUGH+HYDRATE code to assess the production potential of challenging hydrate deposits, i.e., deposits that are characterized by any combination of the following factors: absence of confining boundaries, high thermodynamic stability, low temperatures, low formation permeability. Using high-resolution grids, we show that a new horizontal well design using thermal stimulation coupled with mild depressurization yields production rates that appear modest and insufficient for commercially viable production levels. The use of parallel horizontal wells (with the lower one providing thermal stimulation through heat addition, direct injection or circulation of warm water, and the upper one producing under a mild depressurization regime) offers tantalizing possibilities, and has the potential of allowing commercial production from a very large number of hydrate deposits that are not currently considered as production candidates if the problem of the corresponding large water production can be solved.

### INTRODUCTION

Gas hydrates are solid crystalline compounds in which gas molecules are lodged within a clathrate crystal lattice (Sloan and Koh, 2008). Vast amounts of  $\text{CH}_4$  stored in hydrates in geologic media in the permafrost and in the oceans. The current study is part of a larger effort to determine the technical feasibility of gas production from a wide range of hydrate deposits in geologic media.

Recent studies have determined the conditions, methods and characteristics that enhance production from such deposits. The most important features (Moridis et al., 2008) include (a) high temperatures and pressures (the deepest, warmest deposits are the most desirable), (b) thermodynamic proximity to the H-V-Lw equilibrium conditions (Figure 1), (b) the use of depressurization, because pure thermal stimulation appears to be very slow and ineffective (Moridis and Reagan, 2007a;b), (c) the presence of impermeable boundaries and, in the case of Class 2 systems, thin water zones, and (d) high intrinsic permeabilities of the hydrate-bearing sediments. If these conditions are met, hydrate deposits can yield

methane at high rates (well in excess of 10 MMSCFD) for long periods using conventional production technology (Moridis and Reagan, 2007a;b).

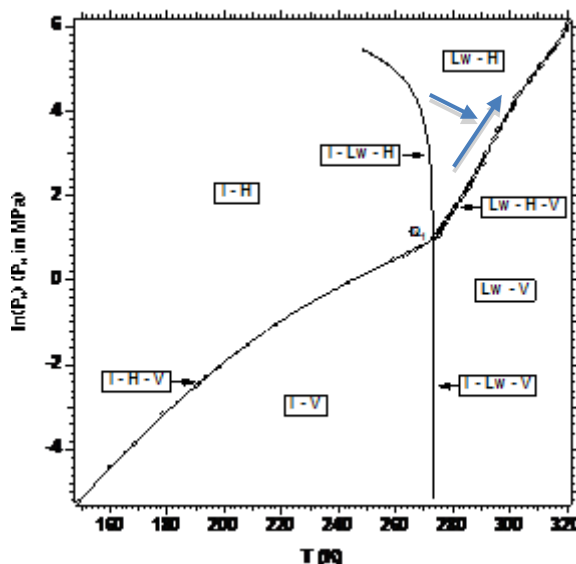


Figure 1. Pressure-temperature equilibrium relationship in the phase diagram of the water- $\text{CH}_4$ -hydrate system (Moridis, 2003). The two arrows show the direction of increasing thermodynamic desirability of a deposit as a production target.

### Challenging Hydrates

In this study we address the issue of gas production from “challenging” hydrates (CG), i.e. those that do not meet the desirability criteria discussed earlier. Such CG include: (a) absence of impermeable boundaries (CG-B), (b) low initial temperatures, and, consequently, pressures (CG-T), (c) increased stability, as indicated from their thermodynamic distance from the hydrate equilibrium conditions (CG-S), (d) extremely low effective permeability  $k_{eff}$ , caused either by very high hydrate saturations  $S_H$  (CG-H), and/or by occurrence in fine-textured sediments, low- $k$  media such as silts and clays (CG-k). An additional type of CG includes hydrate chimneys (CG-C), i.e., marine hydrates that occur at high  $S_H$  in near-vertical cylindrical structures that are associated with past  $\text{CH}_4$  plumes, often extend to the

ocean floor, and usually have limited diameters (usually < 30 m) and no confining boundaries.

Dissociation is orders of magnitude more effective than thermal stimulation as a dissociation method for gas production (Moridis and Reagan, 2007b). However, a common feature of all cases of CG is that depressurization cannot be effectively applied because of (1) the absence of low- $k$  boundaries and high water production (CG-B, CG-C), (2) the very low  $k_{eff}$  (CG-H, CG-k, CG-C), (3) the impracticality of effecting the very large pressure drops needed to cause dissociation of very stable hydrates (CG-S), or (4) low sensible heat to sustainably fuel depressurization-induced dissociation (CG-T). The high cost and progressively diminishing effectiveness of chemical inhibitors precludes their intensive use for gas production from CG, and pure thermal stimulation has been shown to be ineffective (Moridis and Reagan, 2007b). Conjunctive use of thermal stimulation with depressurization appears to be a plausible method for gas production from CGs.

### Objectives

In this study we investigate by means of numerical simulation the production potential of some types of CGs. We focus on CG-B, but we also investigate production from CG-T and CG-k through sensitivity analysis. Additionally, we investigate the effectiveness of two different well designs. We evaluate production according to two criteria: the *absolute* criterion of gas production, and the *relative* criterion of the gas-to-water ratio.

### GEOLOGIC AND NUMERICAL MODEL

The geologic system in this study is based the Tigershark area, located in the Alaminos Canyon Block 818 of the Gulf of Mexico. Log data from a specially designed exploration well in about 2750 m (9000 ft) of water at the site indicated the presence of an 18.25-m (60-ft) thick sandy hydrate-bearing layer (HBL) corresponding to a drilling depth. The HBL has a porosity  $\phi$  of about 0.30 and Darcy-range intrinsic permeability  $k$ . Initial estimates of gas hydrate saturation  $S_H$  derived from analyses of the resistivity and p-wave velocity data indicate a range from 0.6 to over 0.8. Preliminary calculations indicated that the base of the gas hydrate stability zone at this location occurs at or slightly below the base of the HBL. Because of uncertainty about its boundaries, Moridis and Reagan (2007a;b) investigated production from the Tigershark deposit both as a Class 2 deposit (HBL overlying a mobile water zone) and a Class 3 system (HBL bounded by impermeable strata, with no underlying zone of mobile fluids). They showed that the presence of near-impermeable boundaries can yield very high rates (as high as 17 MMSCFD in Class 2, up to 15

MMSCFD in Class 3). However, sensitivity analysis (Reagan et al., 2008; Boswell et al. 2009) indicated that lack of impermeable boundaries can dramatically reduce gas production (Figure 2), while yielding very large amounts of water. Here we investigate the production potential of the Tigershark formation under the hypothesis of a worst-case scenario, i.e., absence of impermeable boundaries, and contact of the hydrate layer with practically infinite aquifers.

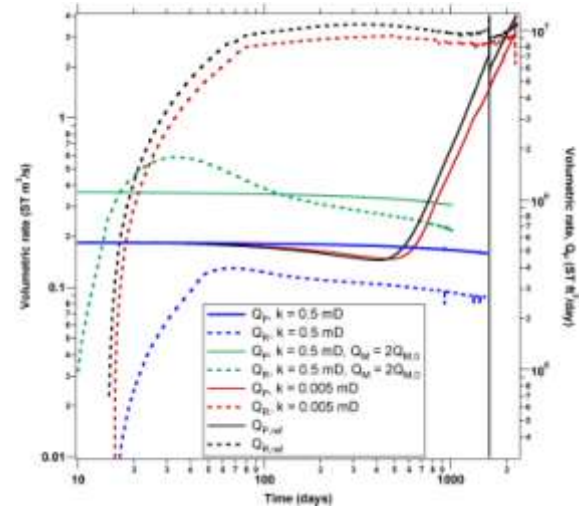


Figure 2. Effect of boundary permeability on gas production from a Tigershark Class 3 deposit (Reagan et al., 2008).

This study involves horizontal wells because of their significant advantages over vertical wells in production from Class 2 and Class 3 deposits (Moridis and Reagan, 2008). We investigated two different horizontal well designs. The single-well of the first design (Figure 3) provides heat to the hydrate-bearing sediment (HBS) by means of hot water that circulates inside the wellbore without coming in contact with the hydrate. The resulting higher  $T$  is expected to promote hydrate dissociation and gas production through the configuration of Figure 3 that operates at a pressure  $P_w$  that is slightly lower than the initial pressure  $P_0$ . By avoiding direct injection of the warm water into the HBS we do not create adverse relative permeability conditions for the flow evolving gas, and the mild depressurization limits the water production. The second design is akin to that used in heavy oil production, and involves two parallel horizontal wells. Heat (through circulation of warm water, electrical or microwave heating, or direct water injection into the HBL) is added to the HBL through the lower well, while the upper well (positioned on the same vertical plane) is the gas collection well operates at a mild depressurization regime. In both well designs, the source of the warm water is assumed to be a deeper warmer reservoir.

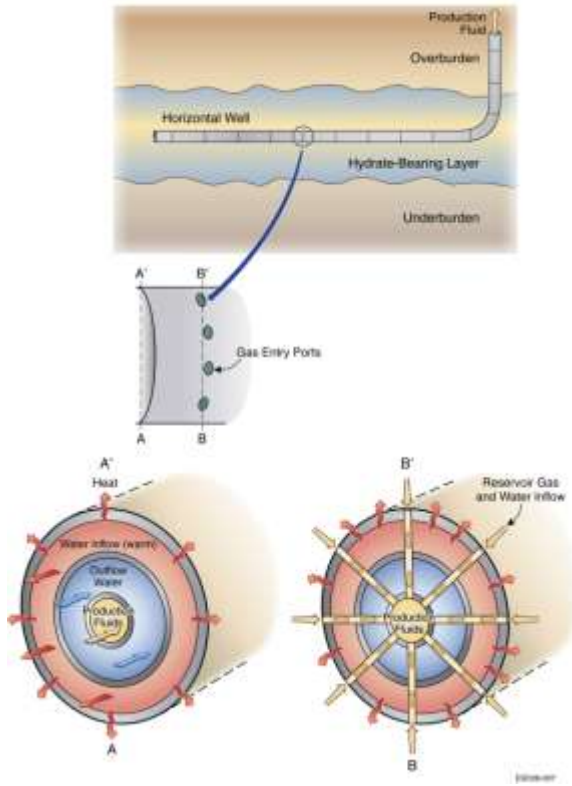


Figure 3. The new well design for concurrent heat addition and gas production.

We conducted the simulations using the TOUGH+HYDRATE code (Moridis et al., 2008). This code can model the non-isothermal hydration reaction, phase behavior, and flow of fluids and heat under conditions typical of natural  $\text{CH}_4$ -hydrate deposits in complex geologic media. It includes both an equilibrium and a kinetic model of hydrate formation and dissociation. The model accounts for heat and up to four mass components (i.e., water,  $\text{CH}_4$ , hydrate, and water-soluble inhibitors such as salts or alcohols) that are partitioned among four possible phases: gas, aqueous liquid, ice, and hydrate. A total of 15 states (phase combinations) can be described by the code, which can handle any combination of hydrate dissociation mechanisms.

We used 2-D grids because of symmetry. The unstructured hybrid grids used in the simulations are shown in Figures 4 and 5, and comprised 47,000 and 27,000 elements, respectively (resulting in 288,000 and 108,000 coupled equations). The system properties and initial conditions are as described by Moridis and Reagan (2007a;b) and shown in Table 1. Both grids had open top and bottom boundaries, i.e., the HBL was connected with the permeable overburden and underburden, allowing fluid and heat flow through the boundaries. The  $x = 40$  m boundary was closed, indicating a well spacing of 80 m.

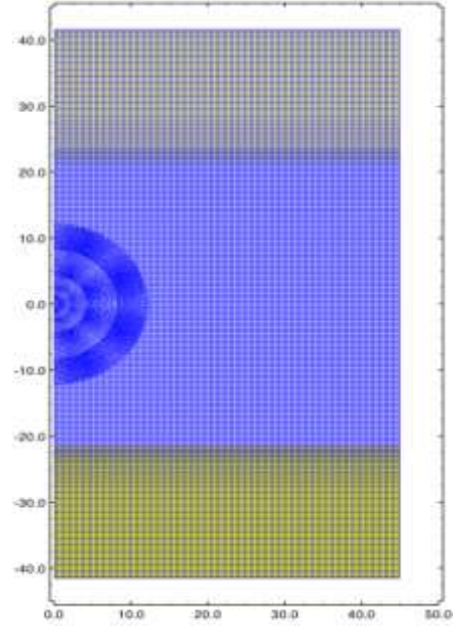


Figure 4. Grid used in the study of the performance of the new well design of Figure 3.

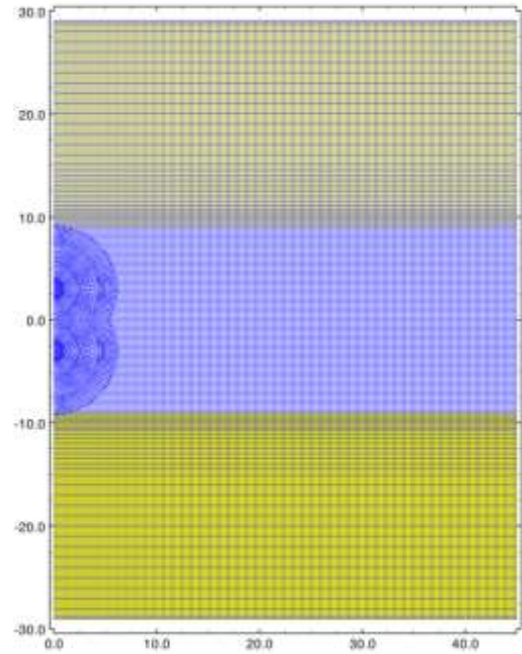


Figure 5. Grid used in the study of the performance of the two parallel horizontal well system.

### THE SINGLE WELL DESIGN

In the evaluation of the single-well design, the reference case involved the sandy HBS ( $k = 7.5 \times 10^{-13} \text{ m}^2$ ) described in the Moridis and Reagan (2007a;b) study, and the temperature of the circulating hot water was  $T_w = 90^\circ \text{C}$ . Sensitivity analysis was investigated through the following additional cases: (a)  $T_w = 120^\circ \text{C}$ , (b) a silty medium with  $k = 7.5 \times 10^{-14}$



m<sup>2</sup>, (c) a clayey medium with  $k = 7.5 \times 10^{-15}$  m<sup>2</sup>, and (d) lower (by 10 °C) initial  $T$ , i.e., a more stable hydrate at the prevailing  $P$ .

Table 1. Physical properties and simulation parameters for the 2-D hydrate-bearing system.

Parameter	Value
Hydrate zone thickness	18.25 m
Initial pressure $P_B$ (at base of HBL)	$3.3 \times 10^7$ Pa
Initial temperature $T_B$ (at base of HBL)	294.15 K (21 °C)
Gas composition	100% CH <sub>4</sub>
Initial saturations in the HBL	$S_H = 0.7, S_A = 0.3$
Water salinity (mass fraction)	0.03
Initial saturations in the HBL	$S_H = 0.7, S_A = 0.3$
Intrinsic permeability $k_r = k_z$ (HBS and boundaries)	$7.5 \times 10^{-13}$ m <sup>2</sup> (= 0.75 D)
Grain density $\rho_R$ (all formations)	2750 kg/m <sup>3</sup>
Dry thermal conductivity	0.5 W/m/K
$k_{ORD}$ (all formations)	
Wet thermal conductivity	3.1 W/m/K
$k_{ORW}$ (all formations)	
Composite thermal conductivity model (Moridis et al., 2005)	$k_{QC} = k_{ORD}$ $+(S_A^{1/2} + S_H^{1/2})(k_{ORW} - k_{ORD})$ $+ f S_I k_{OI}$
Capillary pressure model (vanGenuchten, 1980)	$P_{cap} = -P_0 \left[ (S^*)^{-1/\lambda} - 1 \right]^\lambda$ $S^* = \frac{(S_A - S_{irA})}{(S_{mxA} - S_{irA})}$
$S_{irA}$	1
$l$	0.45
$P_0$	$10^5$ Pa
Relative permeability Model (Moridis et al., 2008)	$k_{rA} = (S_A^*)^n$ $k_{rG} = (S_G^*)^n$ $S_A^* = (S_A - S_{irA}) / (1 - S_{irA})$ $S_G^* = (S_G - S_{irG}) / (1 - S_{irA})$ OPM model
$n$ (from Moridis and Reagan, 2007a;b)	3.572
$S_{irG}$	0.02

### The Reference Case and the $T_w = 120$ °C Case

Figure 6 to 9 show respectively the evolutions of the following variables over time:  $P$ ,  $T$ ,  $S_H$  and  $S_G$ . The very low pressure drop  $\Delta P$  (Figure 6) is evident in that it creates an anomaly fully confined in a limited zone around the well. This is caused by the low effective permeability  $k_{eff}$  of the HBL in the area surrounding the dissociated zone (Figure 8). As expected, the temperature disturbance does not propagate far from the well because of the limited efficiency of conduction as the main heat transfer

mechanism, and the rate of its propagation declines significantly over time as the volume around the well increases as a function of  $r^2$ . A direct consequence of the limited advance of the temperature front is the limited extent of the dissociated region (Figure 8).

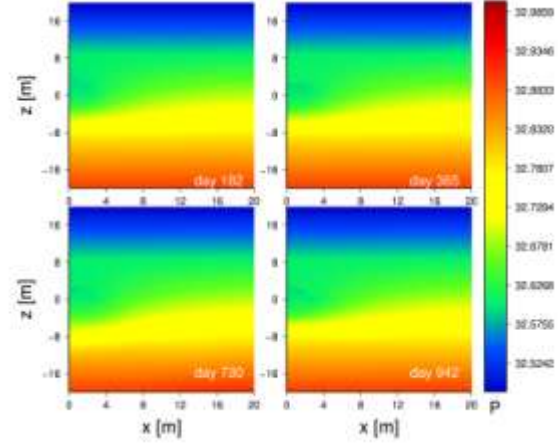


Figure 6. Evolution of pressure  $P$  over time during production from the single well of Figure 3.

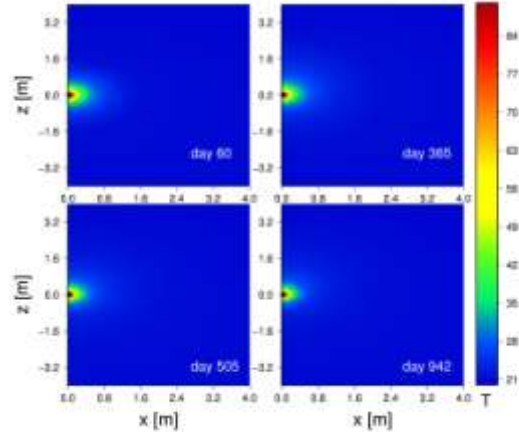


Figure 7. Evolution of temperature  $T$  during production from the single well of Figure 3.

Of particular interest is the high- $S_H$  region immediately ahead of the dissociation front (Figure 8). This occurs because the edge of this front is the locus of local maximum of  $P$  in the system, with fluids moving both away and toward the well. Gas moving deeper into the hydrate body (away from the well) encounters conditions that are conducive to secondary hydrate formation that result in  $S_H$  higher than the initial one. The  $S_G$  distribution in Figure 10 indicates that practically all the dissociated gas that has not been produced is trapped within the hydrate-free cylindrical zone defined by the dissociation front. Because of buoyancy, gas accumulates at the top of the cylindrical dissociated zone, while the water released from dissociation drains and accumulates at the bottom of the cylindrical hydrate-free zone.

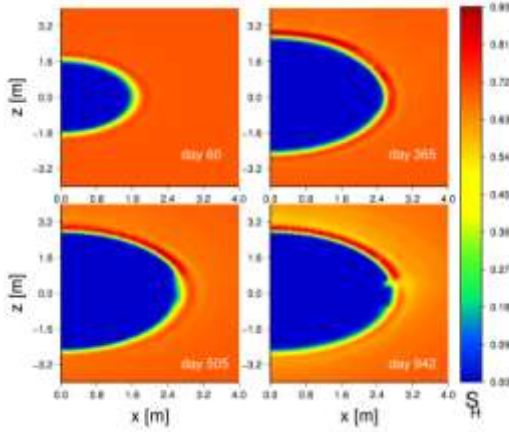


Figure 8. Evolution of hydrate saturation  $S_H$  during production from the single well of Figure 3.

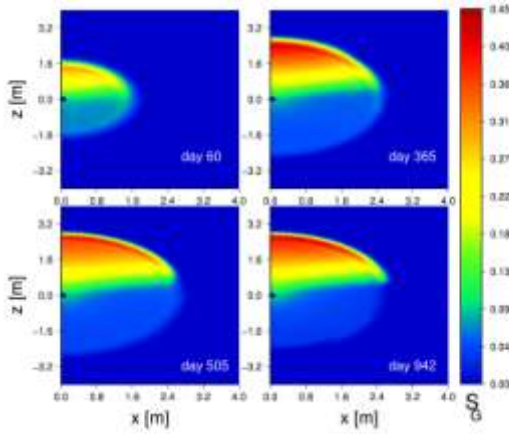


Figure 9. Evolution of gas saturation  $S_G$  during production from the single well of Figure 3.

Figure 10 shows the volumetric rates (per linear m of the horizontal well) of  $\text{CH}_4$  (a) release  $Q_R$ , (b) production in the gas phase  $Q_{PG}$ , and (c) total gas production  $Q_{PT}$ , i.e., both in the gas and aqueous phase for the  $T_w = 90^\circ\text{C}$  and  $120^\circ\text{C}$  cases. Note that  $Q_{PT}$  exceeds  $Q_R$ , and that the majority of the produced gas comes from  $\text{CH}_4$  dissolved in the water rather than from the free gas phase, and that. The production rates appear to be quite low, even if we assume that all the dissolved  $\text{CH}_4$  (a very significant fraction of  $Q_{PT}$ ) is recovered. Additionally, the higher  $T_w$  appears to have a limited effect on gas production, increasing  $Q_{PT}$  only slightly over the  $90^\circ\text{C}$  case. The corresponding water production rates  $Q_w$  and gas-to-water ratios  $R_{GW} = V_P/M_w$  in Figure 12 show the larger  $T_w$  has slight (if any) practical effect, that water production is manageable and that the  $R_{GW}$  is not prohibitively low. However, for a 1000 m well, long-term  $Q_{PT} < 7,000 \text{ ST m}^3$  (245 MSCFD), and about 40 times lower than the rule-of-thumb for commercially viable production rates from offshore gas wells.

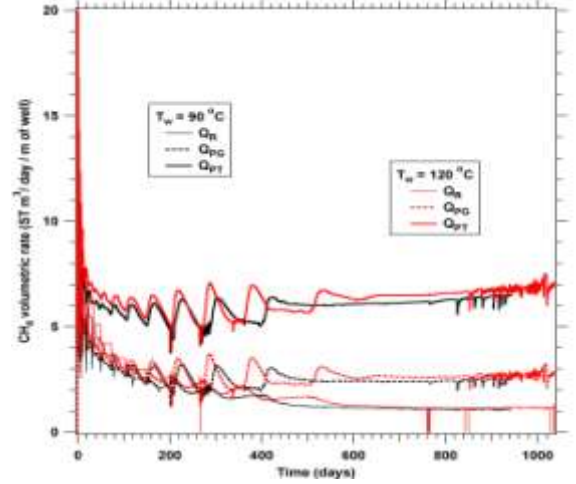


Figure 10. Evolution  $Q_R$ ,  $Q_{PG}$  and  $Q_{PT}$  during production from the single well of Figure 3.

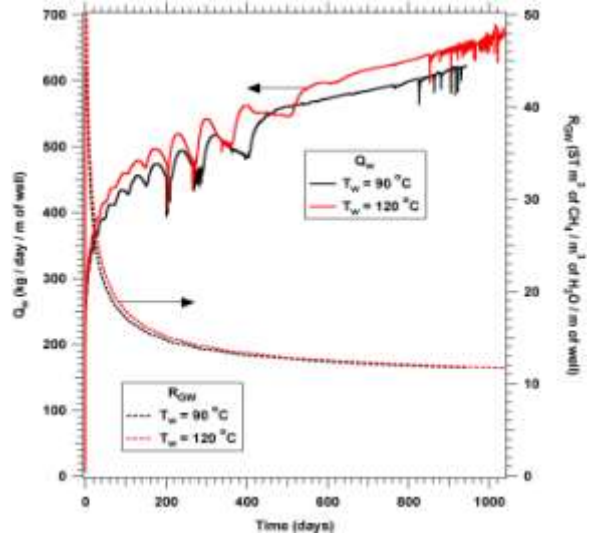


Figure 11. Evolution  $Q_w$  and  $R_{GW}$  during production from the single well of Figure 3

### Effect of Finer Texture-Media (Lower $k$ )

Figure 12 shows  $Q_R$ ,  $Q_{PG}$ , and  $Q_{PT}$  in (a) the reference case of a sandy HBL, (b) the case of a silt, (c) a clay with a low well pressure drop  $\Delta P_w = 1 \text{ atm}$  (Case Clay-LP), and (d) a clay case of a higher  $\Delta P_w = 5 \text{ atm}$  (case Clay-HP). The decreasing permeability  $k$  and increasing capillary pressure  $P_{cap}$  of a progressively finer texture (moving from a sand to a clay) leads to an increasing  $Q_R$ , but a decreasing  $Q_{PT}$ . Thus,  $Q_{PT}$  from the sandy system (already quite low) is the highest of all cases. Additionally, the contribution of production in the gas phase  $Q_{PG}$  increases with a decreasing permeability. As expected, the higher  $\Delta P_w$  leads to higher  $Q_R$ ,  $Q_{PG}$ ,  $Q_{PT}$ ,  $Q_w$ , and  $R_{GW}$ , but  $Q_{PT}$  is still lower than that for the sandy HBS.

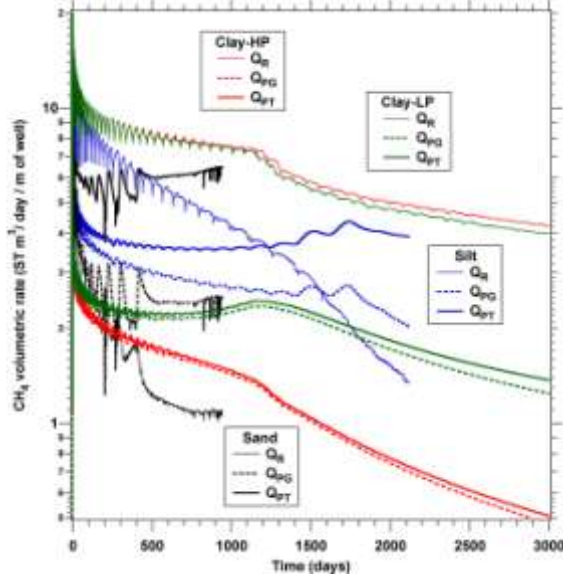


Figure 12. Effect of HBS texture on  $Q_R$ ,  $Q_{PG}$  and  $Q_{PT}$  during production from the single well.

The corresponding  $Q_w$  and  $R_{GW}$  in Figure 13 show that water production increases with the fineness (decreasing  $k$  and increasing  $P_{cap}$ ) of the HBS texture, but  $R_{GW}$  exhibits the opposite pattern. The high  $R_{GW}$  in silt and clay systems (relative criterion) cannot compensate for the low production (absolute criterion).

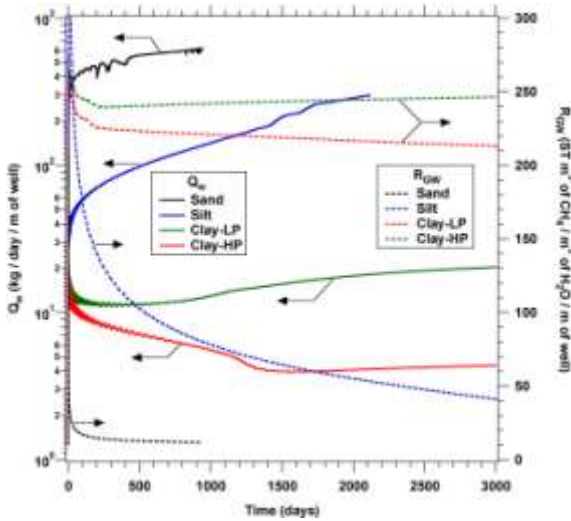


Figure 13. Effect of HBS texture on  $Q_w$  and  $R_{GW}$  during production from the single well.

### Effect of Temperature

Figure 14 shows  $Q_R$ ,  $Q_{PG}$ , and  $Q_{PT}$  in a silty HBL with (a) the reference initial  $T_0 = 21^\circ\text{C}$ , and (b) in a colder system with  $T_0 = 11^\circ\text{C}$  at the same pressure, and indicates that lower  $T_0$  results in significantly lower gas production. Additionally, Figure 15 indicates that the lower  $Q_{PT}$  of the colder case is

further burdened by a lower  $R_{GW}$  despite a decreasing  $Q_w$ . This was expected because the  $k_{eff}$  and the corresponding  $Q_w$  remain low during the longer time it takes for the colder (and thermodynamically more stable) HBL to reach the dissociation temperature at the prevailing pressure.

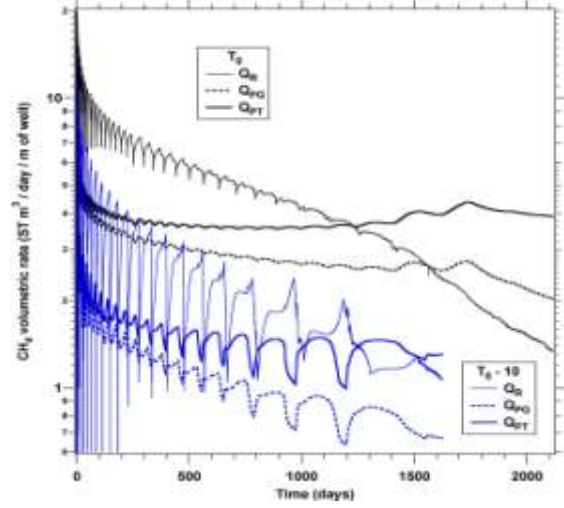


Figure 14. Effect of temperature on  $Q_R$ ,  $Q_{PG}$  and  $Q_{PT}$  during production from the single well.

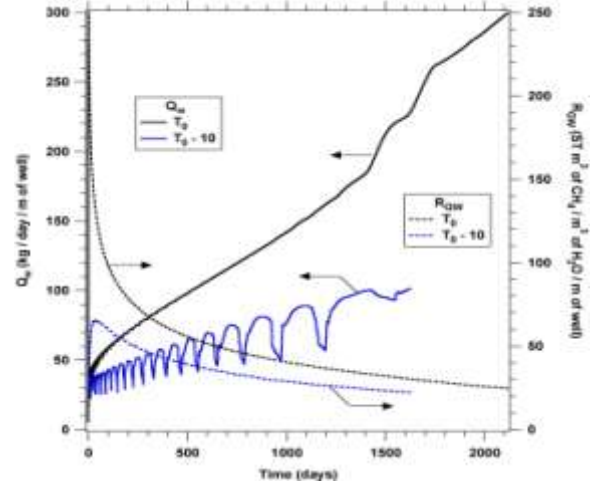


Figure 15. Effect of temperature on  $Q_w$  and  $R_{GW}$  during production from the single well.

### THE TWO-WELL DESIGN

We study the following cases, all involving sandy systems: (a) Case A, with hot water ( $T_w = 90^\circ\text{C}$ ) circulating in the lower well (LW) without entering the HBL, and the upper well (UL) operating at a  $\Delta P_w = 1$  atm, (b) Case B1, with heat added to the HBL at a rate of  $1000\text{ W/m}$ , and a  $\Delta P_w = 1$  atm at the UW, (c) Case B2, differing from B1 in that  $\Delta P_w = 0.2 P_0$ , (d) Case C1, with warm water ( $T_{iw} = 60^\circ\text{C}$ ) injected into the HBL through the LW at a rate of  $Q_{iw} = 5 \times 10^{-3}\text{ kg/s/m}$  of the well, and  $\Delta P_w = 1$  atm at the UW, (e)



Case C2, differing from C1 in that  $\Delta P_w = 0.1 P_0$ , and (f) Case C3, differing from C1 in that  $\Delta P_w = 0.2 P_0$ .

### Cases A, B1 and B2

Figure 16 shows that in Cases A and B1, no  $\text{CH}_4$  is ever produced in the gas phase ( $Q_{PG} = 0$ ), i.e., all the produced gas originated from  $\text{CH}_4$  dissolved in the aqueous phase. The low  $Q_{PT}$  ceases completely after only about 52 days.  $Q_R$  continues past that time because of the continuous heat addition, but this does not lead to continuous gas production because the released gas remains trapped in a dissociated, hydrate-free cylindrical zone that is surrounded by hydrate at very high saturations that exceed the initial  $S_H$  (Figure 17). This is caused by gas from dissociation moving into the HBL and creating secondary hydrates that reach levels resulting in a reduction of  $k_{eff}$  to practically zero. Note that secondary hydrates are also formed around the producing UW. In Figure 17, the gas and water accumulation at the top and the bottom, respectively, of the isolated cylindrical zone are evident.

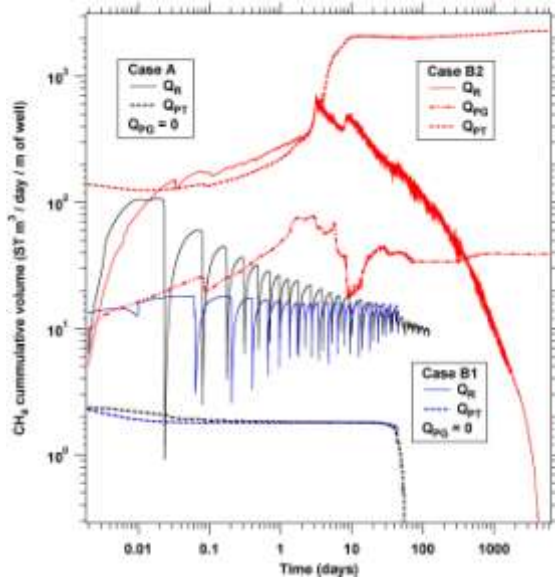


Figure 16. Evolution of  $Q_R$ ,  $Q_{PG}$  and  $Q_{PT}$  during production from the two-well system (Cases A, B1, and B2).

Figure 18 confirms the creation of the isolated zone by showing  $Q_w$  declining to zero at the same time that  $Q_{PT}$  tends to zero. This was expected because no free gas is ever produced, and flow of the aqueous phase is necessary to obtain the dissolved  $\text{CH}_4$  (the only gas source in these cases).

Case B2 appears very different. Figure 16 indicates order of magnitude higher  $Q_R$  and  $Q_{PT}$  in addition to a large  $Q_{PG}$ , all of which are attributable to the stronger depressurization (as the thermal regime remains the same). Thus, the thermal

stimulation in the LW serves only to develop an initial  $Q_R$  reaches a maximum after about 3 days, and then declines continuously until it is reduced to zero after about 4,000 days when the hydrate is exhausted.  $Q_{PG}$  is practically constant after about 60 days, while  $Q_{PT}$  exhibits a jump at 3 days (corresponding roughly to the depressurization front reaching the upper boundary of the HBL), and remains practically constant after 10 days.

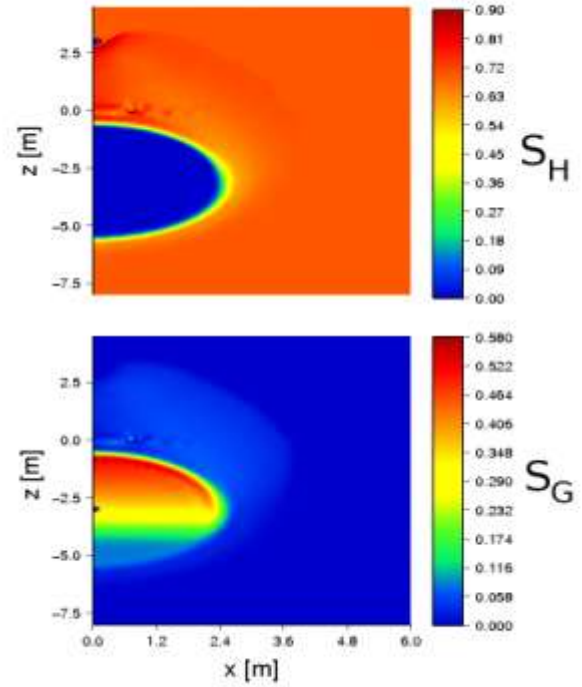


Figure 17. Distribution of  $S_H$  and  $S_G$  during production from the two-well system at  $t = 60$  days (Case A).

This is confirmed by Figure 18, which shows a step increase in  $Q_w$  at  $t = 3$  days, and a constant  $Q_w$  after about 10 days that corresponds to a significant reduction in  $R_{GW}$ . Thus,  $Q_{PT}$  for a 1000 m well system reaches a long-term near-constant level of  $2.2 \times 10^6$  ST  $\text{m}^3/\text{day}$  (76 MMSCFD), this is hampered by a large water production.

The reason for this promising  $Q_{PT}$  performance is the effectiveness of depressurization as a dissociation method, as demonstrated by the  $S_H$  distribution over time in Figure 19. The corresponding  $P$  and  $T$  distributions in Figures 20 and 21, respectively, show the establishment of a steady state pressure regime (because of the permeable boundaries) and the negligible effect of the warm water injection, which appears to be completely overwhelmed by the effects of depressurization (evident by the changes in the  $T$

distribution over time in the vicinity of, and within, the hydrate body).

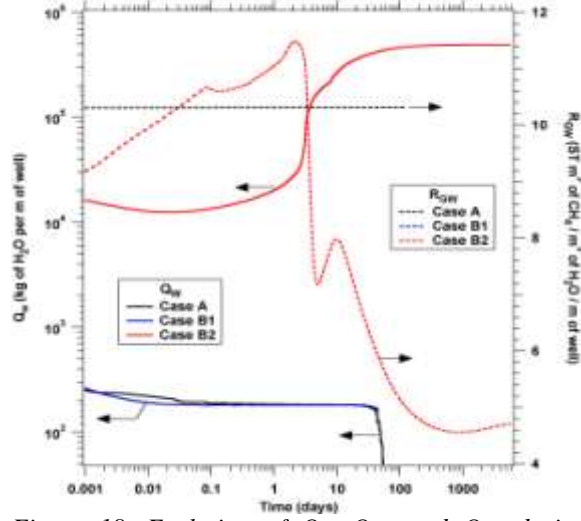


Figure 18. Evolution of  $Q_R$ ,  $Q_{PG}$  and  $Q_{PT}$  during production from the two-well system (Cases A, B1, and B2).

### Cases C1, C2 and C3

The  $Q_R$ ,  $Q_{PG}$ , and  $Q_{PT}$  in Figure 22 indicate that, in Case C1, warm water injection appears to have a worse overall effect than the heat addition methods in Cases A and B1. As in these cases, no  $CH_4$  is ever produced in the gas phase ( $Q_{PG} = 0$ ), but the low  $Q_{PT}$  ceases completely after 21.6 days because of significant secondary hydrate creation around the UW and around the hydrate-free cylindrical zone of complete dissociation (Figure 23), which brings  $k_{eff}$  down to zero levels at these locations. The  $Q_R$  curve is entirely analogous to that of Case A in Figure 16, as are the  $S_H$  and  $S_G$  distributions in Figure 20.

The  $Q_R$ ,  $Q_{PG}$ , and  $Q_{PT}$  of Case C3 in Figure 22 are indistinguishable from those in Case B2, indicating (a) large production potential for this approach, and (b) that the heat addition method plays a minimal role in the pattern of the system response (and possibly only to create an initial high- $k_{eff}$  zone to allow further dissociation by means of depressurization). Figure 24 shows the similarity of the  $Q_W$  and  $R_{GW}$  between Cases A, B1 and C1, in addition to the practical coincidence of the system behavior in Cases B2 and C3.

Case C2 appears quite different. While its  $Q_R$ ,  $Q_{PG}$ , and  $Q_{PT}$  curves in Figure 19 initially appear to track those for Case C2 (albeit at a lower level because of the lower  $\Delta P_w$ ),  $Q_R$  and  $Q_{PG}$  are reduced to zero levels after 34 and 18 days, respectively, because  $\Delta P_w$  is insufficient to prevent the formation of an isolated hydrate-free zone surrounded by high  $S_H$ . Water production continues past this point because the flow to the UW is not blocked (Figure

24), but  $Q_{PT}$  is reduced because dissolved gas is the only source of  $CH_4$ .

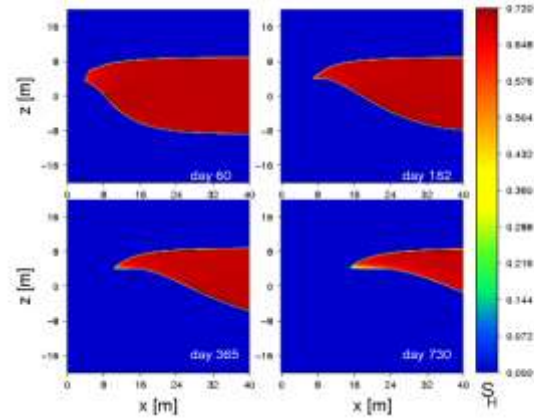


Figure 19. Evolution of hydrate saturation  $S_H$  distribution during production in Case C3.

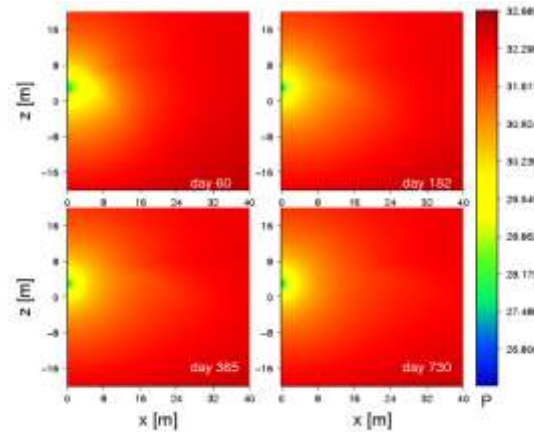


Figure 20. Evolution of pressure  $P$  distribution during production in Case C3.

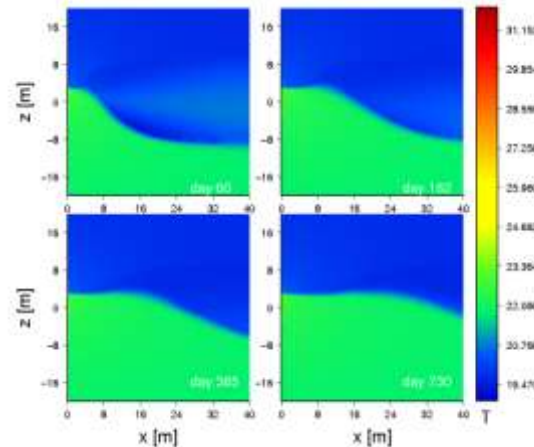


Figure 21. Evolution of reservoir temperature  $T$  distribution during production in Case C3.



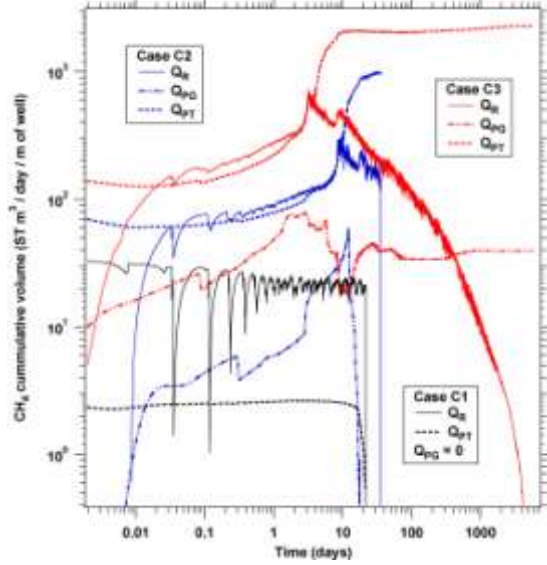


Figure 22. Evolution of  $Q_R$ ,  $Q_{PG}$  and  $Q_{PT}$  during production from the two-well system (Cases C1, C2, and C3).

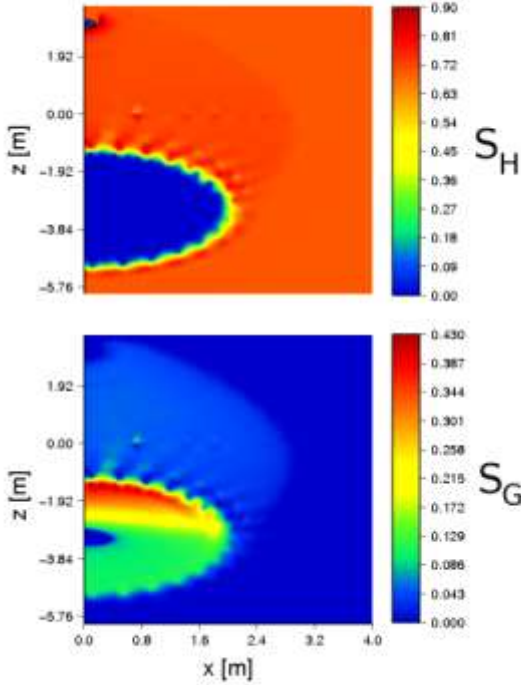


Figure 23. Distribution of  $S_H$  and  $S_G$  during production from the two-well system at  $t = 21.6$  days (Case C1).

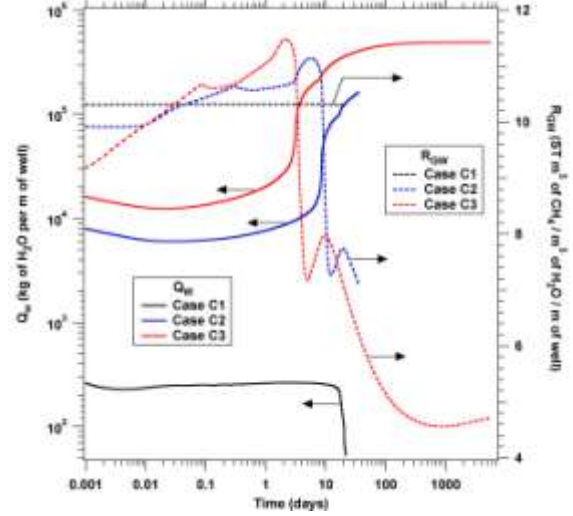


Figure 24. Evolution of  $Q_R$ ,  $Q_{PG}$  and  $Q_{PT}$  during production from the two-well system (Cases C1, C2 and C3).

## CONCLUSIONS

We reach the following conclusions from this study:

- The use of the new, single-well design involving concurrent HBL heating and production from different segments along the same wellbore does not appear a promising solution to the problem of production from challenging hydrates because it results in very low gas production rates and unfavorable gas-to-water ratios.
- Using the new, single-well design, production increases with the coarseness of the HBS (i.e., with an increasing  $k$  and  $P_{cap}$ , thus favoring sandy over silty and clayey HBS) and with the initial temperature of the deposit. Using a lower well (bottomhole) pressure  $P_w$  increases gas production, but it also increases the undesirable water production.
- If the  $P_w$  of the UW is maintained at levels very close to  $P_0$ , the two parallel horizontal well system appears to be very ineffective, resulting in very short production times before flow to the UW is blocked by secondary hydrate that brings the  $k_{eff}$  to zero levels.
- If  $P_w = 0.8 P_0$  in the UW, then depressurization is by far the dominant process and effective dissociation occurs, resulting in high gas production rates (up to 76 MMSCFD). However, the gas production is accompanied by a large water production.

## ACKNOWLEDGMENT

This work was supported by the Assistant Secretary for Fossil Energy, Office of Natural Gas and Petroleum Technology, through the National Energy Technology Laboratory, under the U.S. Department of Energy, Contract No. DE-AC02-05CH11231. The

authors are indebted to John Apps and Dan Hawkes for their careful review.

## **REFERENCES**

- Boswell, R., D. Shelander, M. Lee, T. Latham, T. Collett, G. Guerin, G. Moridis, M. Reagan and D. Goldberg, Occurrence of gas hydrate in Oligocene Frio sand: Alaminos Canyon Block 818: Northern Gulf of Mexico, *J. Mar. Pet. Geo.*, 26, 8, 1499-1512, 2009.
- Moridis, G.J., Numerical Studies of Gas Production from Methane Hydrates, *SPE Journal*, 32(8), 2003.
- Moridis, G.J., Seol, Y., and T. Kneafsey, Studies of reaction kinetics of methane hydrate dissociation in porous media (Paper 1004), *Proceedings of the 5th International Conference on Gas Hydrates*, Trondheim, Norway, 12-16 June 2005.
- Moridis, G.J., and M.T. Reagan, Strategies for Gas Production From Oceanic Class 3 Hydrate Accumulations, OTC-18865, *Proceedings of the 2007 Offshore Technology Conference*, Houston, Texas, 30 April – 3 May, 2007a
- Moridis, G.J., and M.T. Reagan, Gas Production From Oceanic Class 2 Hydrate Accumulations, OTC 18866, *Proceedings of the 2007 Offshore Technology Conference*, Houston, Texas, U.S.A., 30 April–3 May 2007.
- Moridis, G.J., Kowalsky, M.B., and K. Pruess, *TOUGH+HYDRATE v1.0 User's Manual: A Code for the Simulation of System Behavior in Hydrate-Bearing Geologic Media*, Report LBNL-0149E, Lawrence Berkeley National Laboratory, Berkeley, CA.
- Reagan, M.T., Moridis, G.J., and Zhang, K., Sensitivity Analysis of Gas Production from Class 2 and Class 3 Hydrate Deposits, OTC 19554, *Proceedings of the 2008 Offshore Technology Conference*, Houston, Texas, USA, 5-8 May 2008.
- Sloan, E.D. and C. Koh, *Clathrate Hydrates of Natural Gases*, CRC Press, New York, NY, 2008.
- Van Genuchten, M.T., A closed-form equation for predicting the hydraulic conductivity of unsaturated soils, *Soil Sci. Soc.*, 44, 892-898, 1980.